PH575 Spring 2019

World Metrology Day! 20 May 2019

http://www.worldmetrologyday.org/



- SI units underlie global science, trade
- Connected to fundamental constants





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Lecture #21 Semiconductors: optical properties: Kittel Ch. 8 pp. 187-191; Ch 15 pp. 435-444





Figure VI-1-1: Different types of optical absorption phenomena; (1) transitions of highlying bands, (2) excitons, (3) fundamental absorption (VB-to-CB transition and Urbachtail), (4) impurity absorption, (5) free-carrier absorption and (6) Reststrahlen absorption.



Figure VI-1-1: Different types of optical absorption phenomena; (1) transitions of highlying bands, (2) excitons, (3) fundamental absorption (VB-to-CB transition and Urbachtail), (4) impurity absorption, (5) free-carrier absorption and (6) Reststrahlen absorption.



Figure VI-3-1: The absorption edge of (a) a direct semiconductor and (b) indirect semiconductor. The energy E_{vert} marks the threshold for the vertical transition.



Parabolic bands:

$$\begin{split} \hbar \omega_{photon} &= E_g + \frac{\hbar^2 k^2}{2m_e^*} + \frac{\hbar^2 k^2}{2m_h^*} \\ \hbar \omega_{photon} &= E_g + \frac{\hbar^2 k^2}{2\mu}; \quad \frac{1}{\mu} = \frac{1}{m_e^*} + \frac{1}{m_h^*} \end{split}$$

Joint density of states: $D(E) = 0; \quad \hbar \omega < E_g$ $D(E) = \frac{1}{2\pi^2} \left(\frac{2\mu}{\hbar^2}\right)^{3/2} \left(\hbar \omega - E_g\right)^{1/2}; \quad \hbar \omega > E_g$

Figure VI-2-1: Direct absorption in a semiconductor.

Matrix element:



Calculating band gap from absorption and transmission spectra:

In very clean samples at very low temperatures, onset of absorption is sharp and BG identification is trivial, but under other conditions, extraction is harder.

It can be shown (*e.g.* Cardona and Yu) that the imaginary part of the dielectric constant (proportional to absorption coefficient α , remember) for a DIRECT GAP semiconductor is







http://kottan-labs.bgsu.edu/teaching/workshop2001/chapter6.htm

<u>(Johnson [1967])</u>, <u>(Gobeli and Fan [1956])</u>

Calculating band gap from absorption and transmission spectra:

Close to the gap energy, $E - E_q$ is small, and $E \approx E_q$ so

$$\alpha^2 \propto \left(E - E_g\right)$$

A plot of the square of the absorption coefficient against the photon energy is a straight line, and intersects the energy axis at $E = E_g$

Some real data follows, and illustrates that the real world is not like the text book at all!

Calculating band gap from absorption and transmission spectra:

It can be shown (e.g. Cardona and Yu) that the imaginary part of the dielectric constant (proportional to absorption coefficient a, remember) for a INDIRECT GAP semiconductor is

$$\varepsilon''(\omega) \propto \alpha \propto \left(E \pm E_{phonon} - E_g\right)^2$$

$$\alpha^{1/2} \propto \left(E + E_{ph} - E_g \right)$$

A plot of the square root of the absorption coefficient against the photon energy is a straight line, and intersects the energy axis at $E = E_g - E_{ph}$



Figure VI-2-6: The left diagram shows the perturbation of the band edges by Coulomb interaction with inhomogeneously distributed impurities. This leads to the formation of tails of states shown on the right side. The dashed lines show the distribution of states in the unperturbed case.

Mobility Edges/ Minimum Metallic Conductivity



Effect of Disorder: "Urbach Tails"

States near band edge \Rightarrow extra absorption:

$$\alpha \propto \exp\left[-\frac{\left(E_g - E\right)}{k_B T}\right]$$

Exponential, not power law!

Band gap absorption is important for

Optical detection devices. Light with energy greater than the band gap energy excites electrons and holes in a semiconductor. By application of a voltage, the charged electrons and holes can be collected and measured. Light sensor! (This is the basis of a solar cell; not only is light sensed - it is used to generate electricity.)

Optical emission: The reverse process occurs too. Excited electrons and holes recombine, emitting light, primarily at the energy corresponding to the gap. Light emitter! Non radiative recombination processes (phonon emission for example), can hinder the process.



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Absorption coeff. vs. photon energy at different doping levels, *n*-lnSb, T = 130K: **1**. $6.6 \cdot 10^{13} \text{ cm}^{-3}$; **2**. $7.5 \cdot 10^{17} \text{ cm}^{-3}$; **3**. $2.6 \cdot 10^{18} \text{ cm}^{-3}$; **4**. $6 \cdot 10^{18} \text{ cm}^{-3}$; <u>(Ukhanov [1977])</u>.



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Will Excitonic Circuits Change Our Lives? http://opfocus.org/index.php?topic=story&v=2&s=3

18/8/2008

The direct use of light, both for communication and computation, could speed things up. "Our gallium arsenide transistors process signals using indirect excitons instead of electrons," Butov explains. "They are controlled by gate electrodes exactly like electrons in silicon transistors, but, unlike electrons, they are directly coupled with photons, thus bridging the gap between processing and communication." While computation itself may not be faster than electron-based circuits, signal transmission to other devices or parts of the same chip connected by an optical link will definitely be.



Excitonic absorption in BaCuChF (Ch= S, Se, Te)





XRD shows epitaxial quality of BaCuSeF and BaCuTeF films. BaCuSF is polycrystalline



http://kottan-labs.bgsu.edu/teaching/workshop2001/chapter6.htm



$$E_n = -\frac{\mu_{red}}{m_0} \frac{1}{\varepsilon_r^2} \frac{R_H}{n^2} = \frac{R_X}{n^2}$$

$$r_n = \frac{m_0}{\mu_{red}} \varepsilon_r n^2 a_H = n^2 a_X$$

http://kottan-labs.bgsu.edu/teaching/workshop2001/chapter6.htm



We have performed investigation of optical nonlinearities induced by exciton or carrier photogeneration in InGaAs/InP QW structures by using the infrared femtosecond optical parametric amplifier. The absorption spectrum shows clearly visible heavy and light hole exciton peaks



http://www.tf.uni-kiel.de/matwis/amat/semi_en/kap_5/advanced/t5_1_3.html