

Problem 7.42

(a) Faraday's law says $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$, so $\mathbf{E} = 0 \Rightarrow \frac{\partial \mathbf{B}}{\partial t} = 0 \Rightarrow \mathbf{B}(\mathbf{r})$ is independent of t .

(b) Faraday's law in integral form (Eq. 7.18) says $\oint \mathbf{E} \cdot d\mathbf{l} = -d\Phi/dt$. In the wire itself $\mathbf{E} = 0$, so Φ through the loop is constant.

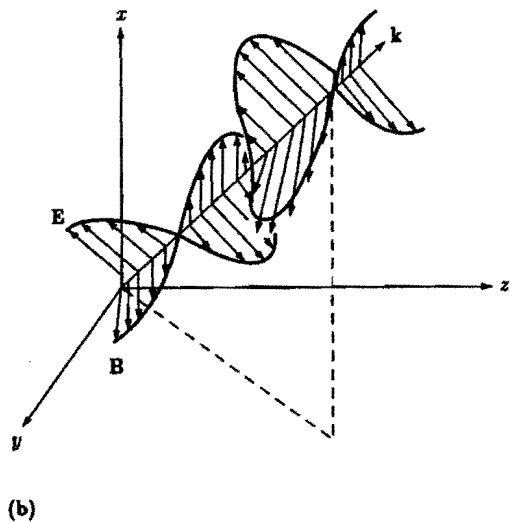
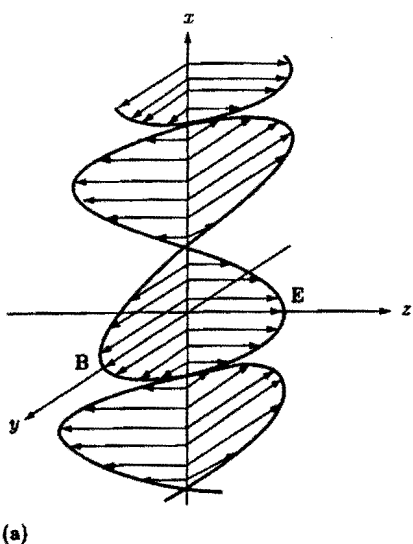
(c) Ampère-Maxwell $\Rightarrow \nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$, so $\mathbf{E} = 0, \mathbf{B} = 0 \Rightarrow \mathbf{J} = 0$, and hence any current must be at the surface.

(d) From Eq. 5.68, a rotating shell produces a uniform magnetic field (inside): $\mathbf{B} = \frac{2}{3} \mu_0 \sigma \omega a \hat{z}$. So to cancel such a field, we need $\sigma \omega a = -\frac{3B_0}{2\mu_0}$. Now $\mathbf{K} = \sigma \mathbf{v} = \sigma \omega a \sin \theta \hat{\phi}$, so $\mathbf{K} = -\frac{3B_0}{2\mu_0} \sin \theta \hat{\phi}$.

Problem 9.9

(a) $\mathbf{k} = -\frac{\omega}{c} \hat{x}; \hat{n} = \hat{z}. \quad \mathbf{k} \cdot \mathbf{r} = \left(-\frac{\omega}{c} \hat{x}\right) \cdot (x \hat{x} + y \hat{y} + z \hat{z}) = -\frac{\omega}{c} x; \quad \mathbf{k} \times \hat{n} = -\hat{x} \times \hat{z} = \hat{y}.$

$\mathbf{E}(x, t) = E_0 \cos\left(\frac{\omega}{c} x + \omega t\right) \hat{z}; \quad \mathbf{B}(x, t) = \frac{E_0}{c} \cos\left(\frac{\omega}{c} x + \omega t\right) \hat{y}.$



(b) $\mathbf{k} = \frac{\omega}{c} \left(\frac{\hat{x} + \hat{y} + \hat{z}}{\sqrt{3}}\right); \hat{n} = \frac{\hat{x} - \hat{z}}{\sqrt{2}}.$ (Since \hat{n} is parallel to the xz plane, it must have the form $\alpha \hat{x} + \beta \hat{z}$;

since $\hat{n} \cdot \mathbf{k} = 0, \beta = -\alpha$; and since it is a unit vector, $\alpha = 1/\sqrt{2}$.)

$\mathbf{k} \cdot \mathbf{r} = \frac{\omega}{\sqrt{3}c} (\hat{x} + \hat{y} + \hat{z}) \cdot (x \hat{x} + y \hat{y} + z \hat{z}) = \frac{\omega}{\sqrt{3}c} (x + y + z); \quad \mathbf{k} \times \hat{n} = \frac{1}{\sqrt{6}} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 1 & 1 & 1 \\ 1 & 0 & -1 \end{vmatrix} = \frac{1}{\sqrt{6}} (-\hat{x} + 2\hat{y} - \hat{z}).$

$\mathbf{E}(x, y, z, t) = E_0 \cos\left[\frac{\omega}{\sqrt{3}c} (x + y + z) - \omega t\right] \left(\frac{\hat{x} - \hat{z}}{\sqrt{2}}\right);$
 $\mathbf{B}(x, y, z, t) = \frac{E_0}{c} \cos\left[\frac{\omega}{\sqrt{3}c} (x + y + z) - \omega t\right] \left(\frac{-\hat{x} + 2\hat{y} - \hat{z}}{\sqrt{6}}\right).$

Problem 9.10

$P = \frac{I}{c} = \frac{1.3 \times 10^3}{3.0 \times 10^8} = 4.3 \times 10^{-6} \text{ N/m}^2.$ For a perfect reflector the pressure is twice as great:

$8.6 \times 10^{-6} \text{ N/m}^2.$ Atmospheric pressure is $1.03 \times 10^5 \text{ N/m}^2$, so the pressure of light on a reflector is

$(8.6 \times 10^{-6}) / (1.03 \times 10^5) = 8.3 \times 10^{-11} \text{ atmospheres.}$

Problem 9.33

(a) (i) Gauss's law: $\nabla \cdot \mathbf{E} = \frac{1}{r \sin \theta} \frac{\partial E_\phi}{\partial \phi} = 0. \checkmark$

(ii) Faraday's law:

$$\begin{aligned} -\frac{\partial \mathbf{B}}{\partial t} &= \nabla \times \mathbf{E} = \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta E_\phi) \hat{\mathbf{r}} - \frac{1}{r} \frac{\partial}{\partial r} (r E_\phi) \hat{\boldsymbol{\theta}} \\ &= \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left[E_0 \frac{\sin^2 \theta}{r} \left(\cos u - \frac{1}{kr} \sin u \right) \right] \hat{\mathbf{r}} - \frac{1}{r} \frac{\partial}{\partial r} \left[E_0 \sin \theta \left(\cos u - \frac{1}{kr} \sin u \right) \right] \hat{\boldsymbol{\theta}}. \\ \text{But } \frac{\partial}{\partial r} \cos u &= -k \sin u; \quad \frac{\partial}{\partial r} \sin u = k \cos u. \\ &= \frac{1}{r \sin \theta} \frac{E_0}{r} 2 \sin \theta \cos \theta \left(\cos u - \frac{1}{kr} \sin u \right) \hat{\mathbf{r}} - \frac{1}{r} E_0 \sin \theta \left(-k \sin u + \frac{1}{kr^2} \sin u - \frac{1}{r} \cos u \right) \hat{\boldsymbol{\theta}}. \end{aligned}$$

Integrating with respect to t , and noting that $\int \cos u dt = -\frac{1}{\omega} \sin u$ and $\int \sin u dt = \frac{1}{\omega} \cos u$, we obtain

$$\mathbf{B} = \frac{2E_0 \cos \theta}{\omega r^2} \left(\sin u + \frac{1}{kr} \cos u \right) \hat{\mathbf{r}} + \frac{E_0 \sin \theta}{\omega r} \left(-k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) \hat{\boldsymbol{\theta}}.$$

(iii) Divergence of \mathbf{B} :

$$\begin{aligned} \nabla \cdot \mathbf{B} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 B_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta B_\theta) \\ &= \frac{1}{r^2} \frac{\partial}{\partial r} \left[\frac{2E_0 \cos \theta}{\omega} \left(\sin u + \frac{1}{kr} \cos u \right) \right] + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left[\frac{E_0 \sin^2 \theta}{\omega r} \left(-k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) \right] \\ &= \frac{1}{r^2} \frac{2E_0 \cos \theta}{\omega} \left(k \cos u - \frac{1}{kr^2} \cos u - \frac{1}{r} \sin u \right) \\ &\quad + \frac{1}{r \sin \theta} \frac{2E_0 \sin \theta \cos \theta}{\omega r} \left(-k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) \end{aligned}$$

$$= \frac{2E_0 \cos \theta}{\omega r^2} \left(k \cos u - \frac{1}{kr^2} \cos u - \frac{1}{r} \sin u - k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) = 0. \checkmark$$

(iv) *Ampère/Maxwell:*

$$\begin{aligned} \nabla \times \mathbf{B} &= \frac{1}{r} \left[\frac{\partial}{\partial r} (r B_\theta) - \frac{\partial B_r}{\partial \theta} \right] \hat{\phi} \\ &= \frac{1}{r} \left\{ \frac{\partial}{\partial r} \left[\frac{E_0 \sin \theta}{\omega} \left(-k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) \right] - \frac{\partial}{\partial \theta} \left[\frac{2E_0 \cos \theta}{\omega r^2} \left(\sin u + \frac{1}{kr} \cos u \right) \right] \right\} \hat{\phi} \\ &= \frac{E_0 \sin \theta}{\omega r} \left(k^2 \sin u - \frac{2}{kr^3} \cos u - \frac{1}{r^2} \sin u - \frac{1}{r^2} \sin u + \frac{k}{r} \cos u + \frac{2}{r^2} \sin u + \frac{2}{kr^3} \cos u \right) \hat{\phi} \\ &= \frac{k E_0 \sin \theta}{\omega r} \left(k \sin u + \frac{1}{r} \cos u \right) \hat{\phi} = \frac{1}{c} \frac{E_0 \sin \theta}{r} \left(k \sin u + \frac{1}{r} \cos u \right) \hat{\phi}. \\ \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} &= \frac{1}{c^2} \frac{E_0 \sin \theta}{r} \left(\omega \sin u + \frac{\omega}{kr} \cos u \right) \hat{\phi} = \frac{1}{c^2} \frac{\omega E_0 \sin \theta}{k r} \left(k \sin u + \frac{1}{r} \cos u \right) \hat{\phi} \\ &= \frac{1}{c} \frac{E_0 \sin \theta}{r} \left(k \sin u + \frac{1}{r} \cos u \right) \hat{\phi} = \nabla \times \mathbf{B}. \checkmark \end{aligned}$$

(b) *Poynting Vector:*

$$\begin{aligned} \mathbf{S} &= \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B}) = \frac{E_0 \sin \theta}{\mu_0 r} \left(\cos u - \frac{1}{kr} \sin u \right) \left[\frac{2E_0 \cos \theta}{\omega r^2} \left(\sin u + \frac{1}{kr} \cos u \right) \hat{\theta} \right. \\ &\quad \left. + \frac{E_0 \sin \theta}{\omega r} \left(-k \cos u + \frac{1}{kr^2} \cos u + \frac{1}{r} \sin u \right) (-\hat{r}) \right] \\ &= \frac{E_0^2 \sin \theta}{\mu_0 \omega r^2} \left\{ \frac{2 \cos \theta}{r} \left[\sin u \cos u + \frac{1}{kr} (\cos^2 u - \sin^2 u) - \frac{1}{k^2 r^2} \sin u \cos u \right] \hat{\theta} \right. \\ &\quad \left. - \sin \theta \left(-k \cos^2 u + \frac{1}{kr^2} \cos^2 u + \frac{1}{r} \sin u \cos u + \frac{1}{r} \sin u \cos u - \frac{1}{k^2 r^3} \sin u \cos u - \frac{1}{kr^2} \sin^2 u \right) \hat{r} \right\} \\ &= \boxed{\frac{E_0^2 \sin \theta}{\mu_0 \omega r^2} \left\{ \frac{2 \cos \theta}{r} \left[\left(1 - \frac{1}{k^2 r^2} \right) \sin u \cos u + \frac{1}{kr} (\cos^2 u - \sin^2 u) \right] \hat{\theta} \right.} \\ &\quad \left. + \sin \theta \left[\left(-\frac{2}{r} + \frac{1}{k^2 r^3} \right) \sin u \cos u + k \cos^2 u + \frac{1}{kr^2} (\sin^2 u - \cos^2 u) \right] \hat{r} \right\}.} \end{aligned}$$

Averaging over a full cycle, using $\langle \sin u \cos u \rangle = 0$, $\langle \sin^2 u \rangle = \langle \cos^2 u \rangle = \frac{1}{2}$, we get the intensity:

$$\mathbf{I} = \langle \mathbf{S} \rangle = \frac{E_0^2 \sin \theta}{\mu_0 \omega r^2} \left(\frac{k}{2} \sin \theta \right) \hat{r} = \boxed{\frac{E_0^2 \sin^2 \theta}{2 \mu_0 c r^2} \hat{r}}.$$

It points in the \hat{r} direction, and falls off as $1/r^2$, as we would expect for a spherical wave.

$$(c) P = \int \mathbf{I} \cdot d\mathbf{a} = \frac{E_0^2}{2 \mu_0 c} \int \frac{\sin^2 \theta}{r^2} r^2 \sin \theta d\theta d\phi = \frac{E_0^2}{2 \mu_0 c} 2\pi \int_0^\pi \sin^3 \theta d\theta = \boxed{\frac{4\pi}{3} \frac{E_0^2}{\mu_0 c}}.$$