

Hydrogen atom

Recall:  $\Delta \Psi(r, \theta, \varphi) + \frac{2\mu}{\hbar^2} (E - V(r)) \Psi(r, \theta, \varphi) = 0$

$\Psi(r, \theta, \varphi) = R(r) Y(\theta, \varphi); \quad Y(\theta, \varphi) \Rightarrow Y_e^m(\theta, \varphi)$

$L^2 Y_e^m(\theta, \varphi) = \hbar^2 l(l+1) Y_e^m(\theta, \varphi); \quad L_z Y_e^m(\theta, \varphi) = \hbar m Y_e^m(\theta, \varphi)$

Now it's time to consider

the radial part  $R(r)$ : from Lecture #2

$\frac{d}{dr} \left( r^2 \frac{dR(r)}{dr} \right) \cdot \frac{1}{R(r)} + \frac{2\mu r^2}{\hbar^2} (E - V(r)) = \lambda$

multiply by  $R(r)$  and  $\frac{-\hbar^2}{2\mu r^2}$ :  $l(l+1)$

$-\frac{\hbar^2}{2\mu r^2} \left( 2r \frac{dR(r)}{dr} + r^2 \frac{d^2 R(r)}{dr^2} \right) + \left( V(r) + \frac{l(l+1)\hbar^2}{2\mu r^2} \right) \cdot R(r) = E R(r)$

$R(r) = E R(r) \quad (5.1)$

$V_{\text{eff}}(r) = V(r) + \frac{l(l+1)\hbar^2}{2\mu r^2}$

↑  
centrifugal potential

↑  
effective potential

Simplify Eq. (5.1) by introducing  $u(r)$  such <sup>(2)</sup>  
 that  $R(r) = \frac{u(r)}{r} \Rightarrow$  then  $\Rightarrow$

$$-\frac{\hbar^2}{2\mu r^2} \left( \frac{2r}{r} \frac{du(r)}{dr} - \frac{2r}{r^2} u(r) + r^2 \frac{d}{dr} \left( \frac{du(r)}{dr} \cdot \frac{1}{r} - \frac{1}{r^2} u(r) \right) \right) + V_{\text{eff}}(r) \frac{u(r)}{r} = E \frac{u(r)}{r}$$

||

$$\frac{d^2 u(r)}{dr^2} \cdot \frac{1}{r} - \frac{du(r)}{dr} \cdot \frac{1}{r^2} + \frac{2}{r^3} u(r) - \frac{1}{r^2} \frac{du(r)}{dr}$$

$$-\frac{\hbar^2}{2\mu r^2} \left( 2 \frac{du(r)}{dr} - \frac{2}{r} u(r) + r \frac{d^2 u(r)}{dr^2} - \frac{du(r)}{dr} + \frac{2}{r} u(r) - \frac{du(r)}{dr} \right) + V_{\text{eff}}(r) \frac{u(r)}{r} = E \frac{u(r)}{r} ; \Rightarrow$$

$$-\frac{\hbar^2}{2\mu r^2} r \frac{d^2 u(r)}{dr^2} + V_{\text{eff}}(r) \frac{u(r)}{r} = E \frac{u(r)}{r} \Rightarrow$$

$$-\frac{\hbar^2}{2\mu} \frac{d^2 u(r)}{dr^2} + V_{\text{eff}}(r) u(r) = E u(r) \quad (5.2)$$

↑  
 looks like a 1D problem,  
 but  $r \geq 0!$  (not  $-\infty < r < \infty$ )

↑ much simpler  
 than (5.1)

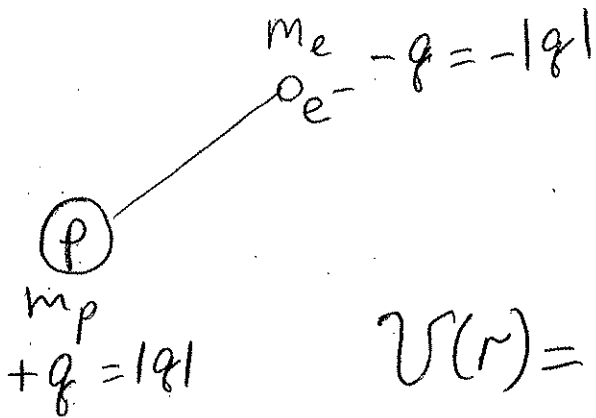
Replace  $u(r) \rightarrow u_{\kappa, l}(r)$   
 $E \rightarrow E_{\kappa, l}$        $\leftarrow V_{\text{eff}}$  depends on  $l$ ,  
 we'll see later that we'll need  $\kappa$

At  $r \rightarrow 0 \Rightarrow R(r) = \frac{u(r)}{r}$  should be finite  $\Rightarrow$  (3)

$$\underline{u_{k,l}(0) = 0}$$

To move any further  $\Rightarrow$  need to specify a potential  $V(r)$  in (5.2) and find  $u_{k,l}(r) \Rightarrow$  the most common example is the hydrogen atom

$\Downarrow$   
proton and electron interact via Coulomb force (electrostatic interaction)



$$V(r) = - \frac{q^2}{4\pi\epsilon_0} \frac{1}{r} = - \frac{e^2}{r}$$

$q = 1.6 \cdot 10^{-19}$  Coulomb      SI system       $\epsilon_0 = 8.85 \cdot 10^{-12} \frac{F}{m}$

$m_e = 9.1 \cdot 10^{-31}$  kg

$m_p = 1.7 \cdot 10^{-27}$  kg

$(m_e \ll m_p) \Rightarrow \mu = \frac{m_e m_p}{m_e + m_p} \approx m_e$

Since the proton is so much heavier, the center of mass almost coincides with the proton, and so the "center-of-mass" particle (recall Lecture #1) is a proton, and the "relative" particle is an electron.

So, now Eq. (5.2) is:

$$(u_{k,l}(0) = 0)$$

$$-\frac{\hbar^2}{2\mu} \frac{d^2 u_{k,l}(r)}{dr^2} + \left( \frac{l(l+1)\hbar^2}{2\mu r^2} - \frac{e^2}{r} \right) u_{k,l}(r) = E_{k,l} u_{k,l}(r) \quad (5.3)$$

First of all  $\Rightarrow$  make all variables dimensionless <sup>(4)</sup>

(recall harmonic oscillator problem)

choose  $\Downarrow$   
 $a_0 = \frac{\hbar^2}{\mu e^2}$  - Bohr radius  
(for  $\text{H}$ -atom  $a_0 \approx 0.5 \text{ \AA}$ )

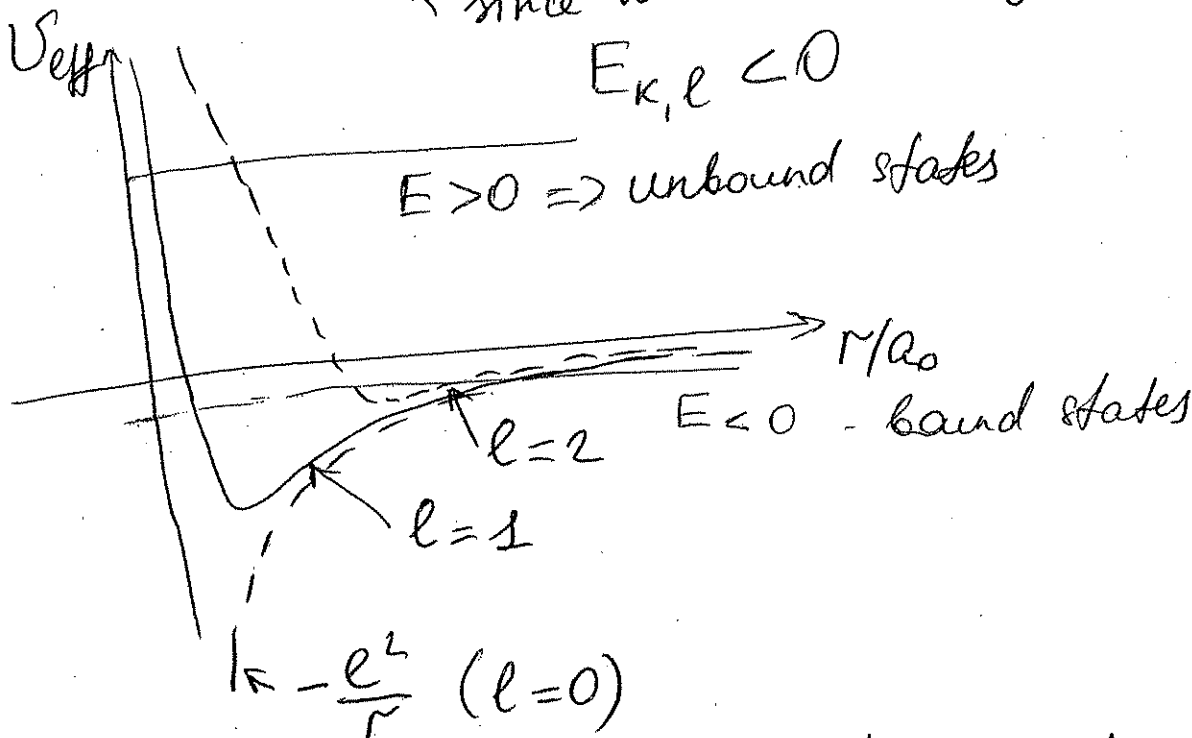
introduce new variables

$$\rho = \frac{r}{a_0}$$

$$E_I = \frac{\mu e^4}{2\hbar^2} \text{ - ionization energy}$$

$$\lambda_{k,l} = \sqrt{-E_{k,l}/E_I}$$

$\nwarrow$  since we are looking for bound states,  
 $E_{k,l} < 0$



With the new variables  $\rho, \lambda_{k,l}$ , the Eq. (5.3) is:

$$\left[ \frac{d^2}{d\rho^2} - \frac{l(l+1)}{\rho^2} + \frac{2}{\rho} - \lambda_{k,l}^2 \right] u_{k,l}(\rho) = 0 \quad (5.4)$$

1) Investigate the asymptotic behavior (5)

$$p \rightarrow \infty : \frac{1}{p^2}, \frac{1}{p} \rightarrow 0 \Rightarrow$$

$$\text{Eq. (5.4) reduces to } \left[ \frac{d}{dp^2} - \lambda_{k,l}^2 \right] u_{k,l}(p) = 0 \Rightarrow$$

$$u_{k,l}(p) = \underbrace{C_1 e^{p\lambda_{k,l}}}_{\rightarrow \infty \text{ as } p \rightarrow \infty} + C_2 e^{-p\lambda_{k,l}}$$

$\rightarrow \infty$  as  $p \rightarrow \infty \Rightarrow$  unphysical  $\Rightarrow$

$$\text{So, at } p \rightarrow \infty \Rightarrow u_{k,l}(p) \sim e^{-p\lambda_{k,l}} \quad C_1 = 0$$

2) Go back to Eq. (5.4) and look for solution in the form  $u_{k,l}(p) = e^{-p\lambda_{k,l}} y_{k,l}(p) \Rightarrow$

$$\left( \frac{d^2}{dp^2} - 2\lambda_{k,l} \frac{d}{dp} + \frac{2}{p} - \frac{l(l+1)}{p^2} \right) y_{k,l}(p) = 0 \quad (5.5)$$

$$\text{Since } u_{k,l}(0) = 0 \Rightarrow y_{k,l}(0) = 0$$

As in the harmonic oscillator problem, look for

$y_{k,l}(p)$  as a power series:

$$y_{k,l}(p) = p^s \sum_{q=0}^{\infty} C_q p^q \quad (C_0 \neq 0) \Rightarrow \text{substitute into (5.5)}$$

$$\frac{dy_{k,l}(p)}{dp} = \sum_{q=0}^{\infty} (q+s) C_q p^{q+s-1}$$

$$\frac{d^2 y_{k,l}(p)}{dp^2} = \sum_{q=0}^{\infty} (q+s)(q+s-1) C_q p^{q+s-2} \quad (6)$$

Then, Eq. (5.5) is:

$$\begin{aligned} \sum_{q=0}^{\infty} (q+s)(q+s-1) C_q p^{q+s-2} & - 2\lambda_{k,l} \sum_{q=0}^{\infty} (q+s) C_q p^{q+s-1} \\ & + 2 \sum_{q=0}^{\infty} C_q p^{q+s-1} - l(l+1) \sum_{q=0}^{\infty} C_q p^{q+s-2} = 0 \end{aligned} \quad (5.6)$$

The lowest power of  $p$  is  $(s-2)$  (in ① and ④ at  $q=0$ )  
collect these terms  $\Rightarrow$

$$(s(s-1)C_0 - l(l+1)C_0) p^{s-2} = 0 \Rightarrow$$

$$\text{since } C_0 \neq 0 \Rightarrow s(s-1) - l(l+1) = 0 \Rightarrow$$

$$s = \begin{cases} l+1 \\ -l \end{cases}$$

$$y_{k,l}(p) = p^{-l} \sum_{q=0}^{\infty} C_q p^q \quad \underline{\text{or}} \quad \boxed{p^{l+1} \sum_{q=0}^{\infty} C_q p^q}$$

$\Downarrow$   
unphysical  
since  $y_{k,l}(0) = 0 \Rightarrow$  discard!

Now change  $q$  to  $q-1$  in ② and ③  $\Rightarrow$

$$\text{②: } \sum_{q=0}^{\infty} (q+s) C_q p^{q+s-1} \Rightarrow \sum_{q=1}^{\infty} (q-1+s) C_{q-1} p^{q+s-2}$$

$$\textcircled{3}: \sum_{q=0}^{\infty} C_q \rho^{q+s-1} \Rightarrow \sum_{q=1}^{\infty} C_{q-1} \rho^{q+s-2} \quad \textcircled{7}$$

Then, Eq. (5.6) is:

$$\sum_{q=1}^{\infty} \left[ ((q+s)(q+s-1) - l(l+1)) C_q + C_{q-1} \cdot (-2\lambda_{k,e} \cdot (q-1+s) + 2) \right] \rho^{q+s-2} = 0 \quad (5.7)$$

(Note that the  $q=0$  term in  $\textcircled{1}$  and  $\textcircled{4}$  was taken care of before, when we found that  $s=l+1$ )

From (5.7)  $\Rightarrow$  obtain the recurrence relation

$$\left[ (q+s)(q+s-1) - l(l+1) \right] C_q = \left[ 2\lambda_{k,e}(q+s-1) - 2 \right] \cdot C_{q-1};$$

Since  $s=l+1 \Rightarrow$

$$\left[ (q+l+1)(q+l) - l(l+1) \right] C_q = \left[ 2\lambda_{k,e}(q+l) - 2 \right] \cdot C_{q-1};$$

$$\begin{aligned} (q+l)^2 + q+l - l^2 - l &= q^2 + 2ql + l^2 + q - l^2 - l \\ &= \underline{q(q+2l+1)} \end{aligned}$$

$$\underline{q(q+2l+1) C_q = 2(\lambda_{k,e}(q+l) - 1) C_{q-1}}$$

If we know  $C_0 \Rightarrow$  can find all other coefficients

Next step: check the convergence of the series

$$\frac{C_q}{C_{q-1}} = \frac{2(\lambda_{k,l}(q+l) - 1)}{q(q+2l+1)} = \frac{2\lambda_{k,l}}{q} \quad (8)$$

If we allow this power series to have infinite number of terms, i.e.  $\sum_{q=0}^{\infty} C_q p^q$ , then this power series would describe an exponential function

$$e^{2p\lambda_{k,l}} = \sum_{q=0}^{\infty} \frac{(2\lambda_{k,l})^q}{q!} p^q, \text{ and so altogether}$$

$$\text{our solution } u_{k,l}(p) = e^{-p\lambda_{k,l}} y_{k,l}(p) = e^{-p\lambda_{k,l}} p^{l+1} e^{2p\lambda_{k,l}} = p^{l+1} e^{p\lambda_{k,l}} \rightarrow \infty$$

need to terminate our power series!

as  $p \rightarrow \infty \Downarrow$   
unphysical

Declare that there exist a  $q_{\max} = k$ ,  
so that  $\lambda_{k,l}(k+l) - 1 = 0$  ( $C_k = 0$ )

Recall  $\rightarrow \lambda_{k,l} = \frac{1}{k+l}$   $\Rightarrow E_{k,l} = -\frac{E_I}{(k+l)^2}$

that this is a dimensionless energy

allowed energy levels ( $k=1, 2, 3, \dots$ )