

Properties of spherical tensors

Example of spherical tensor \Rightarrow spherical harmonics

$$l=0 \Rightarrow Y_0^0 = \frac{1}{\sqrt{4\pi}} \leftarrow \text{const}$$

\uparrow
0th order spherical tensor \rightarrow scalar

$$Y_l^m \Leftrightarrow T_q^{(k)}$$

\uparrow rank
 $\Rightarrow |m| \leq l$

\uparrow component
 $\Rightarrow |q| \leq k$

$$l=1 \Rightarrow Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos\theta, \quad Y_1^{\pm 1} = \mp \sin\theta e^{\pm i\varphi} \sqrt{\frac{3}{8\pi}}$$

$$\begin{cases} x = r \sin\theta \cos\varphi \\ y = r \sin\theta \sin\varphi \\ z = r \cos\theta \end{cases}$$

$$\Rightarrow Y_1^0 = \sqrt{\frac{3}{4\pi}} \frac{z}{r}; \quad Y_1^{\pm 1} = \sqrt{\frac{3}{4\pi}} \frac{1}{r} \left(\mp \frac{x \pm iy}{\sqrt{2}} \right)$$

$$\begin{aligned} \sin\theta e^{\pm i\varphi} &= \sin\theta (\cos\varphi \pm i\sin\varphi) \\ &= \frac{1}{r} (x \pm iy) \end{aligned}$$

Instead of $\vec{r} = (x, y, z)$

Introduce new vector components

$$r_1^1 = -\frac{x+iy}{\sqrt{2}}; \quad r_1^{-1} = \frac{x-iy}{\sqrt{2}}; \quad r_1^0 = z$$

$$\vec{r} = x\vec{i} + y\vec{j} + z\vec{k} = -r_1^1 \vec{e}_1^{-1} - r_1^{-1} \vec{e}_1^1 + r_1^0 \vec{e}_1^0$$

\uparrow
in Cartesian coordinates

\uparrow
in spherical basis

$$\vec{E}_1^1 \stackrel{\leftarrow 0}{=} -\frac{\vec{i} + i\vec{j}}{\sqrt{2}}; \quad \vec{E}_1^0 = \vec{k}, \quad \vec{E}_1^{-1} = \frac{\vec{i} - i\vec{j}}{\sqrt{2}} \quad (2)$$

$$(\vec{E}_1^1 \cdot \vec{E}_1^{-1} = 0)$$

Morning coffee exercise: check that the expansion is

$$\text{So, } r_1^{\pm 1} = \sqrt{\frac{4\pi}{3}} r Y_1^{\pm 1}; \quad r_1^0 = \sqrt{\frac{4\pi}{3}} r Y_1^0 \Rightarrow$$

$$r_1^q \stackrel{\leftarrow \text{rank}}{=} \sqrt{\frac{4\pi}{3}} r Y_1^q \Rightarrow \vec{r} = \sqrt{\frac{4\pi}{3}} r \sum_{q=-1}^1 (-1)^q Y_1^q \vec{E}_1^{-q}$$

↑
expansion in terms of spherical harmonics

Note: In general, any vector \vec{V} can be decomposed in the spherical basis:

$$\vec{V} = V_x \vec{i} + V_y \vec{j} + V_z \vec{k} = \sum_{q=-1}^1 (-1)^q V_q^{(1)} \vec{E}_1^{-q}$$

$$V_{\pm 1}^{(1)} = \mp \frac{V_x \pm i V_y}{\sqrt{2}}; \quad V_0^{(1)} = V_z$$

Also note: Beyond the rank 1, it is no longer possible to express Cartesian and spherical tensors of the same rank in terms of each other.

Recall last lecture: $T_{ij} \Rightarrow T_{ij}^{(0)} + T_{ij}^{(2)} + T_{ij}^{(4)}$

N=2 K=0 T^{(1)K} T^{(2)K}

Cartesian tensor of rank N $\Rightarrow 3^N$ components

Spherical tensor of rank K $\Rightarrow (2K+1)$ -components

$$N=0 \Rightarrow 3^0 = 1 = 2 \cdot 0 + 1; \quad N=1, K=1 \Rightarrow 3^1 = 2 \cdot 1 + 1$$

For any other $N \Rightarrow 3^N \neq 2k+1$

(3)

Selection rules

$$\langle j', m' | T_q^{(k)} | j, m \rangle = 0$$

$$\text{unless } k+j \geq j' \geq |k-j|$$

$$m' = m + q$$

Example

Consider the (H)-atom and its energy levels. In the electric dipole approximation, what are the allowed radiative transitions?

Hint: The probability per unit time for an electric dipole transition from the initial state $|i\rangle$ to the final state $|f\rangle$ is $\sim |M_{fi}|^2$,

where $M_{fi} = \langle f | \vec{E} \cdot (-e\vec{r}) | i \rangle$ is the electric dipole matrix element for this transition. \vec{E} is the polarization of light.

Solution: Since $\vec{E} \cdot (-e\vec{r})$ acts on the orbital angular momentum space and can't change the spin state \Rightarrow use only orbital degrees of freedom.

$$|i\rangle = |n_i, l_i, m_i\rangle, \quad |f\rangle = |n_f, l_f, m_f\rangle$$

$$M_{fi} = -e \langle n_f, l_f, m_f | \vec{E} \cdot \vec{r} | n_i, l_i, m_i \rangle$$

selection rules depend on the polarization of light

Linear polarization: $\vec{E} = (0, 0, 1)$ (along Oz) (4)

$$\vec{E} \cdot \vec{r} = z = r \cos\theta = \sqrt{\frac{4\pi}{3}} r Y_0^0(\theta, \varphi)$$

Circular polarization: $\vec{E} \cdot \vec{r} = \pm \frac{1}{\sqrt{2}} (x \pm iy) = \pm \sqrt{\frac{4\pi}{3}} r Y_{\pm 1}^{\pm 1}(\theta, \varphi)$

$$M_{fi} \sim \int R_{m_f l_f}(r) R_{m_i l_i}(r) r^2 dr$$

for any q

$$q = 0, \pm 1$$

\uparrow \uparrow
 $\parallel z$ circular

$$\int Y_{l_f}^{m_f}(\theta, \varphi) Y_1^q(\theta, \varphi) Y_{l_i}^{m_i}(\theta, \varphi) \sin\theta d\theta d\varphi$$

Use: $Y_l^{\pm m} \sim e^{\pm im\varphi} \Rightarrow$

$$\int_0^{2\pi} e^{-im_f\varphi} e^{iq\varphi} e^{im_i\varphi} d\varphi = \int_0^{2\pi} e^{i(m_i+q-m_f)\varphi} d\varphi =$$

$$= \frac{e^{i(m_i+q-m_f)2\pi} - 1}{i(m_i+q-m_f)} = 0 \text{ unless } \boxed{m_i+q=m_f}$$

$$\Delta m = m_f - m_i = \underset{\substack{\uparrow \\ \text{linear} \\ \parallel Oz}}{0}, \pm 1$$

\uparrow
circular

Are there any selection rules for l ? \Rightarrow (5)

$$\int Y_{l_f}^{m_f*} Y_l^q Y_{l_i}^{m_i} \sin\theta d\theta = a \int Y_{l_f}^{m_f*} Y_{l_i+1}^{m_i+q} \sin\theta d\theta \quad (+)$$

$$\oplus b \int Y_{l_f}^{m_f*} Y_{l_i-1}^{m_i+q} \sin\theta d\theta = \left\{ \begin{array}{l} \text{Properties of } Y_l^m \\ \Downarrow \\ Y_l^q Y_l^m = a Y_{l+1}^{m+q} + b Y_{l-1}^{m+q} \end{array} \right.$$

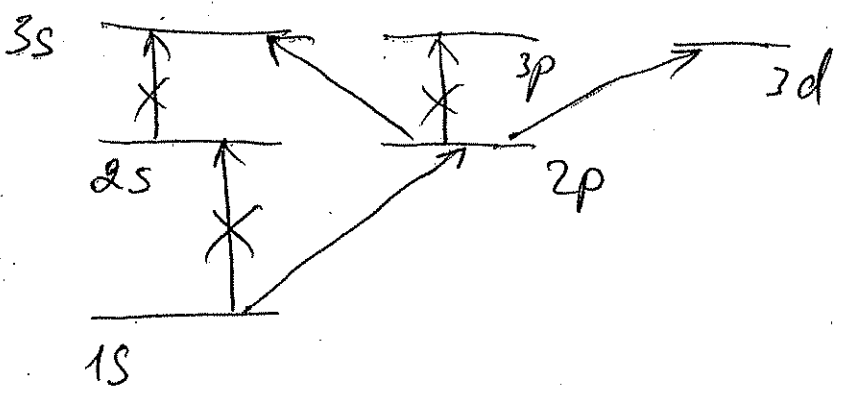
also, see lecture # 17

$$= a \cdot \tilde{C}_1 \delta_{l_f, l_i+1} \delta_{m_f, m_i+q} + b \tilde{C}_2 \delta_{l_f, l_i-1} \delta_{m_f, m_i+q}$$

$$l_f - l_i = \Delta l = \pm 1$$

Note: $\Delta l = 0$ is rejected due to parity

Energy levels of (H)-atom:



$$E_n = -\frac{E_I}{n^2}$$

doesn't depend on l, m

$$\int Y_l^m \sim (-1)^l Y_l^m$$

(in reality \Rightarrow)

$$H = H_{\text{Coulomb}} + H_{\text{SO}} + H_{\text{hf}+-}$$

removes degeneracy)

spin-orbital coupling
spin-spin coupling

