

## Linear spaces

### Definition

A linear (or vector) space is a set of elements (called vectors), with an operation of vector addition and an operation of scalar multiplication.

### Addition rule:

- if  $\Psi$  and  $\Phi$  are vectors (elements) of a space  $\Rightarrow$  their sum is also a vector of the same space
- commutativity:  $\Psi + \Phi = \Phi + \Psi$
- associativity:  $(\Psi + \Phi) + \chi = \Psi + (\Phi + \chi)$
- existence of a zero vector  $0$  such that  $0 + \Psi = \Psi + 0 = \Psi$

- Existence of an inverse vector  $\psi$  such that  $\psi + (-\psi) = (-\psi) + \psi = 0$  (2)

Multiplication rule:

- the product of a scalar with a vector gives another vector. If  $\psi$  and  $\varphi$  are vectors,  $a$  and  $b$  are scalars,  $\Rightarrow a\psi + b\varphi$  is also a vector of the same space.
- Distributivity:  $a(\psi + \varphi) = a\psi + a\varphi$   
 $(a+b)\psi = a\psi + b\psi$
- associativity:  $a(b\psi) = (ab)\psi$
- for each element  $\psi$  there must exist a unitary element  $I$  and a zero scalar such that  $I \cdot \psi = \psi \cdot I = \psi$   
 $0 \cdot \psi = \psi \cdot 0 = 0$

The vector space is called complex or real, depending on whether complex or only real numbers are used as scalars.

## Examples of linear spaces

1) If  ~~$\Psi$~~   $\Psi = (x_1, x_2, \dots, x_n)$

$$\Psi = (y_1, y_2, \dots, y_n)$$

( $x_i, y_i$  - numbers)

with addition defined by

$$\Psi + \Psi = (x_1 + y_1, x_2 + y_2, \dots, x_n + y_n)$$

and multiplication defined by  $a\Psi = (ax_1, ax_2, \dots, ax_n)$

n-dimensional Euclidean space ( $R^n$ )

2) The set of all infinite sequences of numbers

$(x_1, x_2, \dots, x_k, \dots)$  such that  $\sum_{k=1}^{\infty} |x_k|^2$  is finite

with addition and scalar multiplication defined as  
in the previous example  $\Rightarrow l^2$ -space

3) The set of all continuous functions of a real variable  $x$  with an addition of two vectors

$\Psi$  and  $\Psi$  defined by  $(\Psi + \Psi)(x) = \Psi(x) + \Psi(x)$

and multiplication  $(a\psi)(x) = a\psi(x)$

4) The set of all functions  $\psi$  of a real variable  $x$  for which an integral  $\int |\psi(x)|^2 dx$  is finite, with addition and scalar multiplication defined as in the previous example  $\Rightarrow L^2$ -space

What about a "wave function space" (let's call it  $E$ )?  $\Rightarrow E$  is a subspace of  $L^2$ , which in turn is a subspace of Hilbert space.

Formal definition of the Hilbert space:

- it's a linear space also called "inner product"
- there is a so-called scalar product defined as:
  - $(\psi, \varphi) =$  complex number
  - $(\psi, \varphi) = (\varphi, \psi)^*$
  - $(\psi, a\psi_1 + b\psi_2) = a(\psi, \psi_1) + b(\psi, \psi_2)$
  - $(\psi, \psi) = \|\psi\|^2 \geq 0$  (equality only for  $\psi=0$ )

- it's separable } we won't go into heavy
- it's complete } mathematical details  
to see where this comes from

So, each quantum state of a particle is characterized by a state vector  $|\Psi\rangle$ , belonging to an abstract space  $\mathcal{E}$  (which is a subspace of Hilbert space), called the state space of a particle, i.e.  $|\Psi\rangle \in \mathcal{E}$

### Dimension and Basis of a vector space

A set of  $N$  vectors  $\Psi_1, \Psi_2, \dots, \Psi_N$  is said to be linearly independent if and only if the solution of the equation  $\sum_{i=1}^N a_i \Psi_i = 0$  is  $a_1 = a_2 = \dots = a_N = 0$ .

If not all  $a$ 's are zero  $\Rightarrow$  one of the vectors can be expressed as a linear combination of the others  $\Rightarrow \Psi = \sum_{i=1}^N a_i \Psi_i \Rightarrow$  the set  $\{\Psi_i\}$  is linearly dependent.

The dimension of the space is given by ⑥ the maximum number of linearly independent vectors the space can have. If the vectors  $(\psi_1, \psi_2, \dots, \psi_N)$  are linearly independent  $\Rightarrow$  the space is  $N$ -dimensional. In this case, any vector  $\psi$  of the vector space can be expressed as a linear combination

$$\psi = \sum_{i=1}^N a_i \psi_i \quad (1)$$

set of vectors  $\{\psi_i\}$   $\Rightarrow$  base vectors

If the scalar products  $(\psi_i, \psi_j) = \delta_{ij} \Rightarrow$  the basis is orthonormal  $\uparrow$  Kronecker

If the basis spans entire space  $\Rightarrow$  it is complete. The expansion coefficients  $a_i$  in (1) are called the components of the vector  $\psi$  in the basis. Each component  $a_j = (\psi_j, \psi)$ .

(7)

## Examples

1) Are functions  $f(x) = 4, g(x) = x^2, h(x) = e^{3x}$  linearly independent?

Consider  $a_1 f(x) + a_2 g(x) + a_3 h(x) =$   
 $= 4a_1 + a_2 x^2 + a_3 e^{3x} = 0 \Rightarrow$  only if  
 $a_1 = a_2 = a_3 = 0$

$f, g, h$  are linearly independent

2)  $f(x) = x, g(x) = 5x, h(x) = x^2 \Rightarrow ?$   
 linearly independent?

$$a_1 x + a_2 \cdot 5x + a_3 \cdot x^2 = 0$$

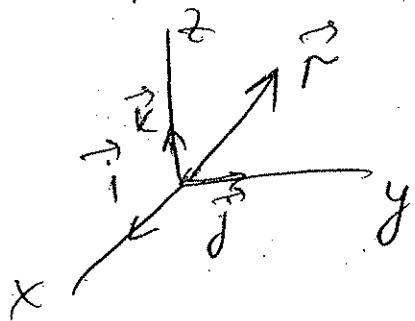
is it valid only if  $a_1 = a_2 = a_3 = 0$ ?

not really  $\Rightarrow x(a_1 + 5a_2 + a_3 x) = 0 \Rightarrow$   
 can be if  $a_3 = 0, a_1 = -5a_2 \neq 0$ !

linearly dependent

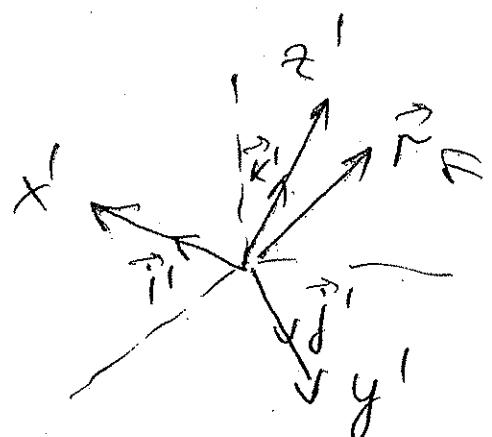
Is the choice of the basis unique?  $\Rightarrow$  no!

Consider



$$\vec{r} = \vec{x}\hat{i} + \vec{y}\hat{j} + \vec{z}\hat{k}$$

expansion of  $\vec{r}$  in  
the  $\vec{i}, \vec{j}, \vec{k}$ -basis



$$\vec{r} = x'\hat{i}' + y'\hat{j}' + z'\hat{k}'$$

note that expansion  
coefficients are  $\left\{ \hat{i}', \hat{j}', \hat{k}' \right.$  - basis  
different in different bases!

Let's go back to wave functions  $\Rightarrow$

we are used to coordinate representation of  
the wave function  $\Rightarrow \Psi(\vec{r}, t)$   $\Rightarrow$  what if  
we don't want to use a coordinate-related  
basis?  $\Rightarrow$  to free state vectors from  
coordinate meaning  $\Rightarrow$  Dirac introduced  
a notation, in which he denoted the state

vector  $\Psi$  by a ket vector  $|\Psi\rangle$  ⑨  
its complex conjugate  $\Psi^*$   $\Rightarrow$  by a bra  
 $\langle \Psi |$ , and the scalar product  $(\Psi, \Psi)$   
(which is one of the attributes of Hilbert  
space, to which state vectors belong!) -

by a bra-ket  $\langle \Psi | \Psi \rangle$

For every ket  $|\Psi\rangle$  there exists a unique bra  
 $\langle \Psi |$  and vice versa. If  $|\Psi\rangle \in E \Rightarrow$   
 $\langle \Psi | \in E^* \in$  dual space of  $E$ .

The ket  $|\Psi\rangle$  represents the system completely,  
i.e. knowing  $|\Psi\rangle$  means knowing all  
amplitudes in all possible representations!!

If we need to find the probability of finding  
the particle at some position in space  $\Rightarrow$

work with the coordinate representation.

The state vector is given by  $\Psi(\vec{r}, t) = \underline{\langle \vec{r}, t | \Psi \rangle}$

In the coordinate representation, the (10)  
scalar product

$$\langle \Psi | \Psi \rangle = \int \Psi^*(\vec{r}, t) \Psi(\vec{r}, t) d^3 r$$

If we are concerned with momentum of  
the particle  $\vec{p} \Rightarrow$  it's more convenient to use  
momentum representation  $\Rightarrow \Psi(\vec{p}, t) = \langle \vec{p}, t | \Psi$

Scalar product  $\Rightarrow \langle \Psi | \Psi \rangle = \int \Psi^*(\vec{p}, t)$ .

•  $\Psi(\vec{p}, t) d^3 p$  momentum space  
wave function

Later we will consider cases when  
it saves a lot of time and effort to  
choose a more suitable representation!

Properties of kets, bras and brackets

• Every ket has a corresponding bra  $\Rightarrow$

$$(|\Psi\rangle)^* = \langle \Psi |$$

$$|a\Psi\rangle = a|\Psi\rangle$$

$$(a|\Psi\rangle)^* = a^* \langle \Psi | ; \quad \langle a\Psi | = a^* \langle \Psi |$$

↑ complex number

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- Scalar product

$$\langle \psi | \psi \rangle + \langle \psi | \psi \rangle \text{ in general case!} \quad \approx$$

$$\underbrace{\langle \psi | \psi \rangle^*}_{\uparrow} = \langle \psi | \psi \rangle$$

$$\begin{aligned}\langle \psi | \psi \rangle^* &= \left( \int \psi^*(\vec{r}, t) \psi(\vec{r}, t) d^3 r \right)^* = \\ &= \int \psi^*(\vec{r}, t) \psi(\vec{r}, t) d^3 r = \langle \psi | \psi \rangle\end{aligned}$$

$$\langle \psi | \psi \rangle = \langle \psi | \psi \rangle \text{ only if } |\psi\rangle, |\psi\rangle \text{ are real!}$$

- The norm  $\langle \psi | \psi \rangle$  is real and positive  
If  $|\psi\rangle$  is normalized  $\Rightarrow \langle \psi | \psi \rangle = 1$

- Schwarz inequality

$$|\langle \psi | \psi \rangle|^2 \leq \langle \psi | \psi \rangle \langle \psi | \psi \rangle$$

if  $|\psi\rangle$  and  $|\psi\rangle$  are linearly dependent.  
(i.e.  $|\psi\rangle = a|\psi\rangle$ )  $\Rightarrow$  equality.

- Triangle inequality

$$\sqrt{\langle \psi + \psi | \psi + \psi \rangle} \leq \sqrt{\langle \psi | \psi \rangle} + \sqrt{\langle \psi | \psi \rangle}$$

if  $|\psi\rangle, |\psi\rangle$  are linearly dependent  $\Rightarrow$  equality

- Orthogonality

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$|\Psi\rangle$  and  $|\Phi\rangle$  are orthogonal if  $\langle \Psi | \Phi \rangle = 0$

- Orthonormality

$|\Psi\rangle$  and  $|\Phi\rangle$  are orthonormal if they are orthogonal and each has a unit norm:

$$\langle \Psi | \Phi \rangle = 0, \quad \langle \Psi | \Psi \rangle = 1, \quad \langle \Phi | \Phi \rangle = 1$$

- Forbidden quantities  $\Rightarrow |\Psi\rangle |\Phi\rangle$  or  $\langle \Psi | \langle \Phi |$

unless they belong to

different spaces, e.g.  $|\Psi\rangle \in E_x$   
 $|\Phi\rangle \in E_y \Rightarrow$

$$|\Psi\rangle \otimes |\Phi\rangle \in E_x \otimes E_y$$

↑

tensor  
product

### Physical meaning of the scalar product

1)  $\langle \Psi | \Phi \rangle$  represents a projection of  $|\Psi\rangle$  onto  $|\Phi\rangle$

2)  $\langle \Psi | \Phi \rangle$  represents the probability amplitude that the system's state  $|\Psi\rangle$  will be found in the state  $|\Phi\rangle$  after a measurement