

PH 585
AMS

Lecture # 10

(1)

Interaction of one-electron atom with external electric & magnetic fields

Static electric fields $\vec{D} \cdot \vec{E}$ ↑ ch. 6
in BSJ

Linear \Downarrow
Stark effect $\Rightarrow H' = eEz$ in $H = H_0 + H'$
electric field $\frac{\hbar^2}{2m} \nabla^2 - \frac{Ze^2}{4\pi\epsilon_0 r}$

Clearly,

$$\Delta E_{1S}^{(1)} = \langle 100 | eEz | 100 \rangle = 0 \leftarrow E \text{ doesn't change energy of } 1S \text{ in the 1st-order approx.}$$

Recall QM:

$n=2$ state $\Rightarrow 2n^2$ -degenerate

Since H' doesn't couple spin states, can consider 4-fold degenerate \Rightarrow

need $\langle \psi_{2\ell m} | H' | \psi_{2\ell m} \rangle \leftarrow$ 4x4 matrix

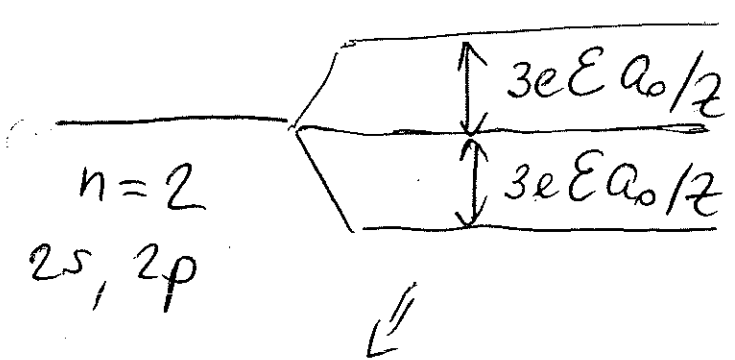
Then, diagonalize the matrix and find energy corrections \Rightarrow

$$\delta \tilde{V} = \pm \frac{3e a_0}{\hbar c} \frac{E}{2} = \pm 12.8 \frac{E}{2} \cdot 10^{-7} \text{ cm}^{-1}$$

$$\Delta E^{(1)} = \pm \frac{3eE a_0}{2} \\ \Delta E^{(1)} = 0$$

$\Delta W = \frac{h}{\tau} \quad 10^{-7} \text{ eV}$

$hc\bar{\nu} = E \quad (2)$



$\frac{1}{\sqrt{2}} (\Psi_{200} - \Psi_{210}) \quad (m=0)$
 $\Psi_{21, \pm 1} \quad (m=\pm 1)$
 $\frac{1}{\sqrt{2}} (\Psi_{200} + \Psi_{210}) \quad (m=0)$

Depending on E and Z could be smaller or larger than, e.g., fine structure splitting \Rightarrow if larger, treat H'_{fine} as a perturbation to $H_0 + H'$; if smaller \rightarrow treat H' as a perturbation to $H_0 + H'_{\text{fine}}$.

Compare: for $n=2 \Rightarrow \bar{\nu}_{\text{fine}} \sim 0.4 \text{ cm}^{-1}$
 (H)-atom
 Lecture # 10

Stark: $\sim 25 E \cdot 10^{-7} \text{ cm}^{-1} = 0.4 \text{ cm}^{-1}$
 $z=1 \quad \frac{12.8 \cdot z}{Z}$

Note: Dielectric strength of air is $\sim 3 \cdot 10^6 \frac{\text{V}}{\text{m}}$
 $\Rightarrow 3 \cdot 10^4 \frac{\text{V}}{\text{cm}}$

Note: $\bar{\nu}_{\text{fine}} \sim Z^4$
 $\bar{\nu}_{\text{Stark}} \sim \frac{1}{Z}$

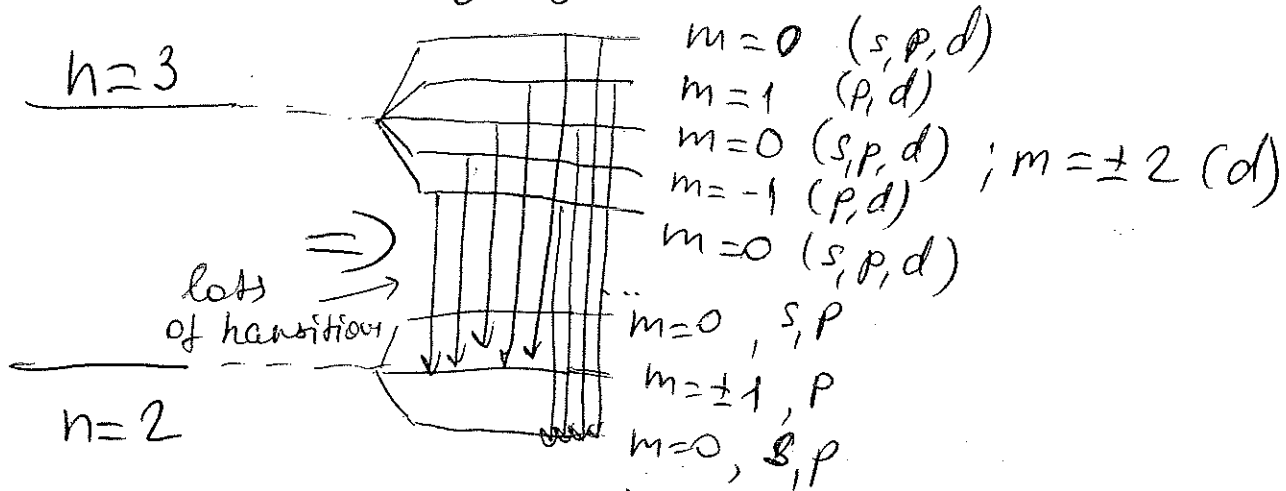
$2^5 = 32$

splitting comparable to fine structure

So, if $Z=2 \Rightarrow$ need $E \sim Z^5 E_{\text{(H)}}$ to get comparable effects

Consider strong fields and H_{α} -line

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H_0

$H_0 + H'$
 eEz

Selection rules:
 (electric dipole approx.)

$\Delta l = ±1$
 $\Delta m = 0, ±1$

\uparrow π -transitions

\nwarrow σ -transitions

Note:

l is not a good quantum number anymore, since states with different l 's are mixed! \Rightarrow e.g. $\frac{1}{\sqrt{2}} (\psi_{2s} \pm \psi_{2p0})$

Recall: $2p \rightarrow 1s \Rightarrow \tau = 1.6 \text{ ns}$
 $2s \rightarrow 1s \Rightarrow \tau = \frac{1}{7} \text{ s}$ } if $E=0$

Now, if $E \neq 0 \Rightarrow 2s$ is mixed with unstable

Consider time evolution of $\psi(\vec{r}, t) = C_1 \psi_{2p0} e^{-\frac{i}{\hbar} E t}$
 superpositional state

$$+ C_2 \Psi_2 e^{-\frac{i}{\hbar} E' t} = \frac{1}{\sqrt{2}} \left[\Psi_1(\vec{r}) e^{-\frac{i}{\hbar} (3e\mathcal{E}a_0)t} + \right. \quad (4)$$

$$E = E_{n=2} - 3e\mathcal{E}a_0 \quad \cos() - i\sin()$$

$$E' = E_{n=2} + 3e\mathcal{E}a_0$$

$$\text{and } \Psi(\vec{r}, t=0) = \Psi_{2s}(\vec{r})$$

$$+ \left. \Psi_2(\vec{r}) e^{\frac{i}{\hbar} (3e\mathcal{E}a_0)t} \right\} e^{-\frac{i}{\hbar} E_{n=2} t} \quad \text{⊖}$$

$\cos() + i\sin()$ ↑

$$t=0 \Rightarrow \frac{1}{\sqrt{2}} (\Psi_1(\vec{r}) + \Psi_2(\vec{r})) = \Psi_{2s}$$

$$\text{⊖} \left(\Psi_{2s} \cos \frac{3e\mathcal{E}a_0 t}{\hbar} + i \Psi_{2p0} \sin \frac{3e\mathcal{E}a_0 t}{\hbar} \right) e^{-\frac{i}{\hbar} E_{n=2} t}$$

$$\Downarrow \quad \frac{1}{\sqrt{2}} (\Psi_2 - \Psi_1)$$

So, the atom oscillates between 2s and 2p0

with a period $T = \frac{\pi \hbar}{3e\mathcal{E}a_0}$

2p → 1s

$$\text{If } \mathcal{E} = 10^7 \frac{\text{V}}{\text{m}} \Rightarrow T = 1.3 \cdot 10^{-12} \text{ s} \ll 1.6 \text{ ns}$$

$$\uparrow$$

$$10^5 \frac{\text{V}}{\text{cm}}$$

can make

2s → 1s happen

so, on average, population in 2s ≈ population in 2p0 during 2p → 1s transition time

Quadratic Stark effect

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For example, $\langle 1s | eEz | 1s \rangle = 0 \Rightarrow$ consider higher perturbation order

$$E_{1s}^{(2)} = e^2 \mathcal{E}^2 \sum_{n \neq 1, \ell, m} \frac{E_{1s}^{(1)}}{E_1 - E_n} |\langle \psi_{n\ell m} | z | \psi_{100} \rangle|^2 = -e^2 \mathcal{E}^2 \sum_{n \neq 1, \ell, m} \frac{|\langle \psi_{n\ell m} | z | \psi_{100} \rangle|^2}{E_n - E_1} > -e^2 \mathcal{E}^2 \frac{1}{E_2 - E_1}$$

$$\sum_{n \neq 1, \ell, m} |\langle \psi_{n\ell m} | z | \psi_{100} \rangle|^2 \quad E_n - E_1 \geq E_2 - E_1$$

$$\underbrace{\sum_{n \neq 1, \ell, m} |\langle \psi_{n\ell m} | z | \psi_{100} \rangle|^2}_{\text{'' } \langle \psi_{100} | z | \psi_{100} \rangle = 0}$$

$$\sum_{n\ell m} |\langle \psi_{n\ell m} | z | \psi_{100} \rangle|^2 = \sum_{n\ell m} \langle \psi_{100} | z | \psi_{n\ell m} \rangle \underbrace{\langle \psi_{n\ell m} | z | \psi_{100} \rangle}_{\text{'' } I}$$

$$\langle \psi_{n\ell m} | z | \psi_{100} \rangle = \langle \psi_{100} | z^2 | \psi_{100} \rangle =$$

$$= \frac{a_0^2}{Z^2} \Rightarrow E_{1s}^{(2)} > -\frac{e^2 \mathcal{E}^2}{\frac{3}{4} E_I Z^2} \frac{a_0^2}{Z^2} = -\frac{8}{3} (4\pi\epsilon_0)$$

$$E_2 - E_1 = -E_I^2 Z^2 \left[\frac{1}{4} - 1 \right] = \frac{3}{4} E_I^2 Z^2 \quad \frac{a_0^3}{Z^4} \mathcal{E}^2$$

Exact evaluation: $E_{1s}^{(2)} = -2.25 (4\pi\epsilon_0) \frac{a_0^3}{z^4} E^2$ (6)

Introduce dipole moment

$$\vec{D} = \alpha \vec{E} = -\frac{\partial E_{1s}^{(2)}}{\partial E}$$

static dipole polarizability of atom in 1s

$$\alpha = 2e^2 \sum_{n \neq 1, l, m} \frac{|\langle \psi_{nlm} | z | \psi_{100} \rangle|^2}{E_n - E_1}$$

Note: selection rules: $\Delta l = \pm 1$
 $\Delta m = 0$

$$E_{1s}^{(2)} = -\frac{1}{2} \alpha E^2 \Rightarrow$$

$$\alpha = 4.5 \cdot (4\pi\epsilon_0) \frac{a_0^3}{z^4}$$

very small!

at $E = 10^8 \frac{V}{m} \Rightarrow \tilde{\nu} = 0.02 \text{ cm}^{-1}$

$10^6 \frac{V}{cm}$ (H)-atom

← have used this already

Rayleigh scattering

⇒ Σ over p-states with $m=0$

Note: if E is very strong \Rightarrow can induce ionisation \Rightarrow e.g. need $\geq 6 \cdot 10^{10} \frac{V}{m}$ to ionise (H)-atom from ground state, but only $\sim 7 \cdot 10^9 \frac{V}{m}$ if the atom is in $n=30$ state

$E_{\text{critical}} \sim \frac{1}{n^4}$